

Indistinguishable particles as their own environment in time-dependent quantum many-body dynamics

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New trends in Quantum Thermodynamics, University of Surrey, July 8, 2024

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Indistinguishable particles as their own environment in timedependent quantum many-body dynamics

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Time-evolution of quantum systems

External perturbation

observed from local observables



The time-dependent Schrödinger equation

$$i\hbar\partial_t |\psi\rangle = \widehat{H}|\psi\rangle$$

$$\begin{split} \hat{H} = & \sum_{j=1}^{N_e} \frac{(\vec{p}_j + e\vec{A}/c)^2}{2m_e} + V_e(\vec{r}_j) + \sum_{i < j} \frac{e^2}{|\vec{r}_i - \vec{r}_j|} & \text{electrons} \\ & + \sum_{n=1}^{N_n} \frac{(\vec{P}_n^2 + Z_n e\vec{A}/c)^2}{2M_n} + V_n(\vec{r}_n) + \sum_{n < m} \frac{e^2}{|\vec{R}_n - \vec{R}_m|} & \text{nuclei} \\ & - \sum_{j,n} \frac{Z_n e^2}{|\vec{r}_j - \vec{R}_n|} & \text{interaction} \end{split}$$

Open quantum systems



Wavefunction replaced by density matrix

$$|\psi\rangle \rightarrow \widehat{\rho_S}$$

Schrödinger equation replaced by Lindblad equation

$$i\hbar\partial_t\hat{\rho}_S = \boxed{[\hat{H}_S,\hat{\rho}_S]} + \underbrace{\sum_j \gamma_j \left(L_j\hat{\rho}_S L_j^{\dagger} - \frac{1}{2} \left\{ L_j^{\dagger} L_j,\hat{\rho}_S \right\} \right)}$$

Decoherence and dissipation

Time-evolution of quantum systems

External perturbation





Outlook:

- Time-dependent two-particle reduced density matrix method (TD2RDM)
- Application to fermionic condensation



1RDM: $D_1 = N \operatorname{Tr}_{2...N} |\Psi\rangle \langle \Psi|$ $D_1(\vec{r}, \vec{r'}, t) = N \int d^3 r_2 \dots d^3 r_N \psi(\vec{r}, \vec{r_2}, \dots, \vec{r_N}) \psi^*(\vec{r'}, \vec{r_2}, \dots, \vec{r_N})$





The physics of the 2RDM

Cumulant decomposition



Closure of equations of motion

Equation of motion for the 2RDM

...



Time-dependent two-particle reduced density matrix (TD2RDM) method

N. N. Bogoliubov and K. P. Gurov. Kinetic Equations in Quantum Mechanics. JETP 17, 614 (1947)

P. C. Martin and J. Schwinger. Theory of Many-Particle Systems. I. Phys. Rev. 115,1342 (1959)

K.-J. Schmitt, P.-G. Reinhard, and C. Toepffer. Truncation of time-dependent manybody theories. Z. Phys. A 336, 123 (1990)

W. Cassing and A. Pfitzner. Self-consistent truncation of the BBGKY hierarchy on the two-body level. Z. Phys. A 342, 161 (1992)

M. Tohyama and P. Schuck. Truncation scheme of time-dependent density-matrix approach. Eur. Phys. J. A 50, 1 (2014)

B. Schäfer-Bung and M. Nest. Correlated dynamics of electrons with reduced twoelectron density matrices. Phys. Rev. A 78, 012512 (2008) A. Akbari et al. PRB 85, 235121 (2012)

Approximate solution of equations of motion

N-representability problem $\Psi \leftarrow D_{12}$

Reconstruction of $D_{123}[D_{12}]$

F. Lackner et al., Phys. Rev. A 91, 023412 (2015)
F. Lackner et al., Phys. Rev. A 95, 033414 (2017)
S. Donsa et al., Phys. Rev. Research 5, 033022 (2023)

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Cumulant decomposition of D_{123}



The three-particle cumulant



This guarantees contraction consistency

$$D_{12} = \frac{1}{N-2} \operatorname{Tr}_3 D_{123}^R$$

F. Lackner et al., Phys. Rev. A 91, 023412 (2015)
F. Lackner et al., Phys. Rev. A 95, 033414 (2017)
S. Donsa et al., Phys. Rev. Research 5, 033022 (2023)

The N-representability problem

Properties of the 2RDM:

- Hermitian
- Positive-semidefinite
- Antisymmetric
- Normalized

Not enough to guarantee

 $D_{12} \leftarrow |\Psi\rangle$

N-representability problem: What are the **necessary and sufficient properties** to guarantee the existence of at least one many-body wavefunction that contracts to the given reduced density matrix?

A. J. Coleman Rev. Mod. Phys. 35, 668 (1963)D. Mazziotti Phys. Rev. Lett 108, 263002 (2012)

Purification

Enforcing **positive semi-definiteness** of 2RDM and two-hole RDM

Space of 2RDMs belonging to a wavefunction (N-representability)

A. J. Coleman Rev. Mod. Phys. 35, 668 (1963)D. Mazziotti Phys. Rev. Lett 108, 263002 (2012)





Outlook:

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- Application to fermionic condensation

System under investigation – large lattice systems

Fermi-Hubbard model in 1D





Correlation dynamics $M_{\rm s} = 20$ U = 0.1 J20 T 2.0 TD2RDM Density fluctuations disappear. 15 --1.5 sites -1.0 **Entropies** -0.5 5 0.0 $D_1 = \sum n_i |n_i\rangle \langle n_i| \implies S_1 = -\sum n_i \ln(n_i)$ 20 60 80 100 120 20-_r2.0 TDHF -1.5 15 - $D_{12} = \sum g_i |g_i\rangle\langle g_i| \implies S_2 = -\sum g_i \ln(g_i)$ sites -1.0 -0.5 5 0.0 20 40 60 80 100 120 0 $t[J^{-1}]$



Fermionic quasi-condensation

Large S₂ goes hand in hand with macroscopic occupation of one eigenstate of the **2RDM**



I. Březinová et al. Phys. Rev. B 109, 174308 (2024)

Fermionic quasi-condensation

Large S₂ goes hand in hand with macroscopic occupation of one eigenstate of the **2RDM**

$$D_{12} = \sum g_i |g_i\rangle\langle g_i|$$
 fermionic condensation if

Increase number of sites up to $M_s = 60$

At half filling this corresponds to a Hilbert space dimension of $\sim 1.4 \times 10^{34}$





 $g_1 \sim O(N)$

I. Březinová et al. Phys. Rev. B 109, 174308 (2024)

Summary and Conclusions

External

- New method based on time-dependent two-particle reduced density matrix (TD2RDM)
- Suitable to describe correlated dynamics of multi-electron systems
- Even strong correlation effects well captured





• Connecting TD2RDM with formalism of open quantum systems



$$i\hbar\partial_t \hat{D}_{12} = [\hat{H}_{12}, \hat{D}_{12}] + \text{Tr}_3[W_{13} + W_{23}, D_{123}] + \sum_j \gamma_j \left(L_j \hat{D}_{12} L_j^{\dagger} - \frac{1}{2} \left\{ L_j^{\dagger} L_j, D_{12} \right\} \right)$$