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Indistinguishable particles as their own environment in time-dependent quantum many-body dynamics

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New trends in Quantum Thermodynamics, University of Surrey, July 8, 2024

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 $W[W|T|F]$

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Indistinguishable particles as their own environment in timedependent quantum many-body dynamics

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Time-evolution of quantum systems

External perturbation

observed from local observables

The time-dependent Schrödinger equation

$$
i\hbar\partial_t|\psi\rangle = \widehat{H}|\psi\rangle
$$

Solution impossible!

$$
\hat{H} = \sum_{j=1}^{N_e} \frac{(\vec{p}_j + e\vec{A}/c)^2}{2m_e} + V_e(\vec{r}_j) + \sum_{i < j} \frac{e^2}{|\vec{r}_i - \vec{r}_j|} \qquad \text{electrons} \n+ \sum_{n=1}^{N_n} \frac{(\vec{P}_n^2 + Z_n e\vec{A}/c)^2}{2M_n} + V_n(\vec{r}_n) + \sum_{n < m} \frac{e^2}{|\vec{R}_n - \vec{R}_m|} \qquad \text{nuclei} \n- \sum_{j,n} \frac{Z_n e^2}{|\vec{r}_j - \vec{R}_n|} \qquad \text{interaction}
$$

Open quantum systems

Wavefunction replaced by density matrix

$$
|\psi\rangle \rightarrow \widehat{\rho_S}
$$

Schrödinger equation replaced by Lindblad equation

$$
i\hbar\partial_t\hat{\rho}_S = \underbrace{\left[\hat{H}_S,\hat{\rho}_S\right]}_{j} + \underbrace{\left\{\sum_j\gamma_j\left(L_j\hat{\rho}_S L_j^\dagger - \frac{1}{2}\left\{L_j^\dagger L_j,\hat{\rho}_S\right\}\right\}\right)}_{j}
$$

Decoherence and dissipation

Time-evolution of quantum systems

External perturbation

Outlook:

- Time-dependent two-particle reduced density matrix method (TD2RDM)
- Application to fermionic condensation

 $D_1 = N \text{Tr}_{2...N} |\Psi\rangle \langle \Psi |$ **1RDM:** $D_1(\vec{r}, \vec{r}', t) = N \int d^3r_2 \dots d^3r_N \, \psi(\vec{r}, \vec{r}_2, \dots, \vec{r}_N) \psi^*(\vec{r}', \vec{r}_2, \dots, \vec{r}_N)$

The physics of the 2RDM

Cumulant decomposition

Closure of equations of motion

Equation of motion for the 2RDM

…

Time-dependent two-particle reduced density matrix (TD2RDM) method

N. N. Bogoliubov and K. P. Gurov. Kinetic Equations in Quantum Mechanics. JETP 17, 614 (1947)

P. C. Martin and J. Schwinger. Theory of Many-Particle Systems. I. Phys. Rev. 115,1342 (1959)

K.-J. Schmitt, P.-G. Reinhard, and C. Toepffer. Truncation of time-dependent manybody theories. Z. Phys. A 336, 123 (1990)

W. Cassing and A. Pfitzner. Self-consistent truncation of the BBGKY hierarchy on the two-body level. Z. Phys. A 342, 161 (1992)

M. Tohyama and P. Schuck. Truncation scheme of time-dependent density-matrix approach. Eur. Phys. J. A 50, 1 (2014)

B. Schäfer-Bung and M. Nest. Correlated dynamics of electrons with reduced twoelectron density matrices. Phys. Rev. A 78, 012512 (2008) A. Akbari et al. PRB 85, 235121 (2012)

Approximate solution of equations of motion

N-representability problem $\Psi \leftarrow D_{12}$

F. Lackner et al., Phys. Rev. A 91, 023412 (2015) F. Lackner et al., Phys. Rev. A 95, 033414 (2017) S. Donsa et al., Phys. Rev. Research **5**, 033022 (2023)

Reconstruction of

 $\overline{D_{123}[D_{12}]}$

Mount Everest [Luca Galuzzi](https://commons.wikimedia.org/wiki/User:Lucag), Wikipedia, [CC BY-](https://creativecommons.org/licenses/by-sa/2.5)SA 2.5

Cumulant decomposition of D_{123}

The three-particle cumulant

This guarantees contraction consistency

$$
D_{12} = \frac{1}{N-2} \text{Tr}_3 D_{123}^R
$$

F. Lackner et al., Phys. Rev. A 91, 023412 (2015) F. Lackner et al., Phys. Rev. A 95, 033414 (2017) S. Donsa et al., Phys. Rev. Research 5, 033022 (2023)

The N-representability problem

Properties of the 2RDM:

- Hermitian
- Positive-semidefinite
- Antisymmetric
- Normalized

Not enough to guarantee

 $D_{12} \leftarrow |\Psi\rangle$

N-representability problem: What are the **necessary and sufficient properties** to guarantee the existence of at least one many-body wavefunction that contracts to the given reduced density matrix?

A. J. Coleman Rev. Mod. Phys. 35, 668 (1963) D. Mazziotti Phys. Rev. Lett 108, 263002 (2012)

Purification

Enforcing **positive semi-definiteness** of 2RDM and two-hole RDM

Space of 2RDMs belonging to a wavefunction (N-representability)

A. J. Coleman Rev. Mod. Phys. 35, 668 (1963) D. Mazziotti Phys. Rev. Lett 108, 263002 (2012)

Outlook:

- Time-dependent two-particle reduced density matrix method D₂RDM)
- Application to fermionic condensation

System under investigation – large lattice systems

Fermi-Hubbard model in 1D

Correlation dynamics $M_{\rm s} = 20$ $U = 0.1 J$ $20 T^2.0$ TD2RDM Density fluctuations disappear. $15 +1.5$ $rac{6}{5}$ 10 - $+1.0$ **Entropies** -0.5 5 -0.0 $D_1 = \sum_{i} n_i |n_i\rangle\langle n_i| \longrightarrow S_1 = -\sum_{i} n_i \ln(n_i)$ 60 80 100 20 120 $20 \sqrt{2.0}$ **TDHF** $+1.5$ $15 D_{12} = \sum g_i |g_i\rangle\langle g_i| \longrightarrow S_2 = -\sum g_i \ln(g_i)$ $\frac{26}{5}$ 10 $\frac{1}{5}$ $+1.0$ ± 0.5 5 0.0 20 40 60 80 100 120 Ω $t [J^{-1}]$

Fermionic quasi-condensation

Large S₂ goes hand in hand with **macroscopic occupation** of one **eigenstate of the 2RDM**

I. Březinová et al. Phys. Rev. B 109, 174308 (2024)

Fermionic quasi-condensation

Large S_2 goes hand in hand with **macroscopic occupation** of one **eigenstate of the 2RDM**

$$
D_{12} = \sum g_i |g_i \rangle \langle g_i|
$$
fermionic condensation if
$$
g_1 \sim O(N)
$$

$$
g_1 \sim O(N)
$$

Increase number of sites up to $M_s = 60$

At half filling this corresponds to a Hilbert space dimension of \sim 1.4 \times 10³⁴

The effect probably disappears in the thermodynamic limit

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Summary and **Conclusions**

External

- New method based on time-dependent two-particle reduced density matrix (TD2RDM)
- Suitable to describe correlated dynamics of multi-electron systems
- Even strong correlation effects well captured

$$
i\hbar\partial_t\hat{D}_{12} = [\hat{H}_{12}, \hat{D}_{12}] + \text{Tr}_3[W_{13} + W_{23}, D_{123}] + \sum_j \gamma_j \left(L_j \hat{D}_{12} L_j^{\dagger} - \frac{1}{2} \left\{ L_j^{\dagger} L_j, D_{12} \right\} \right)
$$